Untangling the complexity of cosmic flows (in galaxy clusters): accretion, mergers, shock waves and turbulence Astrophysics Talks @ INAF/OAS Bologna

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Outlines

Galaxy clusters: the crossroad of Cosmology and Astrophysics

- 2 Cosmological simulations with MASCLET
- The complexity of cosmic flows around galaxy clusters
- 4 The complexity of cosmic flows around cosmic voids

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Galaxy clusters

• The largest and most massive gravitationally bound objects

(Voit, 2005)



 $\begin{array}{l} \mbox{X-ray} \\ \mbox{ICM at} \sim 10^7 {\rm K} \\ \mbox{Thermal bremsstrahlung} \\ \sim 10 \, {\rm keV} \\ \mbox{L}_X \sim 10^{43} - 10^{45} \, {\rm erg/s} \end{array}$

Optical

 $\sim 100-1000$ of member galaxies Lensing of background galaxies

Microwaves

(Sunyaev-Zel'dovich) ICS of CMB photons by high-energy electrons Independent of distance

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Galaxy clusters as cosmological probes

- Almost a century of history... (Zwicky, 1933)
- Cluster counts: extremely sensitive to cosmological parameters



• And much more ...

- Constraining baryon (gas) fraction
- Determination of H₀ and distance measurement using X-ray + tSZ
- tSZ power spectrum (C_{ℓ})
- Determination of the LSS 'peculiar' velocity field (measuring velocities)
- Constraining Dark Energy EoS and its evolution, neutrinos, etc.

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(Allen et al., 2011)

Galaxy clusters as astrophysical laboratories

- The nearly closed-box nature of deep cluster gravitational potentials act as veritable astrophysical laboratories
 - Galaxy formation (Naab & Ostriker, 2017)
 - Effects on the surrounding IGM (e.g., AGN feedback, metal enrichment)



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Understanding galaxy clusters: the Kaiser (1986) self-similar model

Why self-similarity?

- Primordial density fluctuations can be described by a power-law spectrum ⇒ scale-free initial conditions
- Gravity does not have a preferred scale

As long as gravity is the driving mechanism for GC formation, their properties ought to be described by scale-free relations (power laws)

• X-ray scaling relations (in their simplest form; c.f. Giodini, 2013):

$$egin{aligned} M_{\Delta_z} \propto R_{\Delta_z}^3 \ M_{\Delta_z} \propto T_{
m gas}^{3/2} \ L_X \propto T_{
m gas}^2 \end{aligned}$$

 $f_{
m gas}$ independent of M_{Δ_z} $S \propto T_{
m gas} \propto M_{\Delta_z}^{2/3}$

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Understanding galaxy clusters beyond the self-similar model: numerical simulations

Gas (collisional component): hydrodynamics in an expanding frame

$$\frac{\partial \delta}{\partial t} + \frac{1}{a} \nabla \cdot [(1+\delta)\mathbf{v}] = 0 \qquad \qquad \text{EoS: } p = (\gamma - 1)\rho\varepsilon$$

$$\frac{\partial \mathbf{v}}{\partial t} + \frac{1}{a} (\mathbf{v} \cdot \nabla)\mathbf{v} + H\mathbf{v} = -\frac{1}{a} \nabla \phi - \frac{1}{\rho a} \nabla p \qquad \qquad \Rightarrow \text{Can be written as a}$$

$$\frac{\partial E}{\partial t} + \frac{1}{a} \nabla \cdot [(E+p)\mathbf{v}] = -3H(E+p) - H\rho v^2 - \frac{\rho \mathbf{v}}{a} \nabla \phi$$

$$conservation laws$$

DM (collisionless): *N*-Body $\mathbf{v} = a(t) \frac{\mathrm{d}\mathbf{x}}{\mathrm{d}t} \implies 6N$ $\frac{\mathrm{d}\mathbf{v}}{\mathrm{d}t} = -\frac{1}{a(t)} \nabla \phi - H(t) \mathbf{v}$ Poisson equation $\nabla^2 \phi = 4\pi G a^2 \rho_B \delta_T, \quad \delta_T = \delta + \delta_{\mathrm{DM}} + \delta_* + 2$ Additional physics (§3.1.3) Cooling & heating, stars, feedback, etc.

Dolag, 2008; Borgani & Kravtsov, 2011; Kravtsov & Borgani, 2012; Planelles et al., 2015

 $\delta \equiv \frac{\rho - \rho_B}{\rho_B}$: density contrast, **v**: peculiar velocity, ε : specific internal energy, p: pressure, γ : adiabatic index $E \equiv \rho \varepsilon + \frac{1}{2}\rho v^2$: total energy density, ϕ : peculiar gravitational potential, $H(z) \equiv \frac{a}{a}$: Hubble parameter c

Understanding galaxy clusters outskirts

• The exciting new frontier of cluster astrophysics... (Walker, 2019)



Cosmological simulations with MASCLET



(Quilis, 2004; Quilis+2020)

- MASCLET: Mesh Adaptive Scheme for CosmologicaL structurE evoluTion
- Eulerian hydrodynamics based on HRSC techniques for the gas

 N-Body (Particle-Mesh) for evolving DM

- Coupled by gravity solver (Poisson eqn)
- On top of that,
 - Cooling & UV heating
 - Star formation
 - SN and AGN feedback
- Recently ⇒ MHD version (Quilis+20)

- All these elements are built into an AMR scheme:
 - Can be varied upon application
 - Hierarchy of refinement levels with $\Delta x_{\ell+1}/\Delta x_{\ell} = 1/2$



Cosmological simulations with MASCLET: analysis tools

- ASOHF: adaptive SO (DM) halo finder (Planelles+10, Knebe +11)
- HALMA: adaptive FoF (stellar) halo finder (Navarro-González+13)
- Shock Finder (Planelles+13)
- SPEV: radiative transfer code (Mimica+09, Cuesta-Martínez+15, Mimica+16, Planelles+18)
- vortex: Helmholtz-Hodge decomposition in AMR data (+ extensions; Vallés-Pérez+21a,b)
- Void finder (Ricciardelli+13,14, Vallés-Pérez+21c)

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Galaxy cluster mergers and accretion of matter

- How do galaxy clusters grow?
 - Which are the main channels for clusters mass growth?
 - Do DM and baryons behave differently?

$$\Gamma_{\Delta} = \frac{\mathrm{d}\log M_{\Delta}}{\mathrm{d}\log a}$$
(Diemer & Kravtsov, 2014)

- Most of the mass growth occurs during mergers
- However, a large amount of gas is '*left behind*' and slowly reaccreted
- Larger differences between DM and gas for less massive systems



Vallés-Pérez, Planelles & Quilis (2020, MNRAS 499, 2303)

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Accretion of gas onto galaxy clusters

• Does gas fall isotropically into the cluster, or does it fall through some preferential directions?

$$i_M = rac{\Delta M}{R_{
m vir}^2 \Delta \Omega \Delta t}$$

- How does this evolve with cosmic time?
- Does the gas from all directions present the same properties (e.g., entropy, temperature, baryon fraction, etc.)



Vallés-Pérez, Planelles & Quilis (2020, MNRAS 499, 2303)

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Turbulence in galaxy clusters (I: numerical tools)

Helmholtz-Hodge decomposition (Vallés-Pérez, Planelles & Quilis, 2021, CPC 263, 107892)

$$\mathbf{v} = \mathbf{v}_{c} + \mathbf{v}_{s}$$
$$\nabla \times \mathbf{v}_{c} = \mathbf{0} \implies \mathbf{v}_{c} = -\nabla\phi$$
$$\nabla \cdot \mathbf{v}_{s} = \mathbf{0} \implies \mathbf{v}_{s} = \nabla \times \mathbf{A}$$

Had been done for fixed grids (e.g., Vazza+17), but never for AMR data.

It is easy to show that ϕ , **A** are the solution to:

$$abla^2 \phi = -
abla \cdot \mathbf{v} \qquad
abla^2 \mathbf{A} = -
abla imes \mathbf{v}$$

Elliptic PDEs, formally equivalent to Poisson equations!

- FFT (Green's Method) for the base grid (if periodic boundaries)
- Iterative methods (SOR) for the AMR grids



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Hydrodynamical processes in galaxy clusters

Turbulence in galaxy clusters (I: numerical tools)

Reynolds decomposition

(Vallés-Pérez, Planelles & Quilis, 2021, MNRAS 504, 510)

$$\mathbf{v}(\mathbf{x}) \equiv \langle \mathbf{v} \rangle(\mathbf{x}) + \delta \mathbf{v}(\mathbf{x})$$

 $\langle \mathbf{v} \rangle(\mathbf{x}) = \iiint_V \mathrm{d} \mathbf{y} \, \mathbf{v}(\mathbf{x} + \mathbf{y}) G_L(\mathbf{y}; \mathbf{x})$

ICM turbulence is extremely *multi-scale*: $L = L(\mathbf{x})$.

Again, there were algorithms for fixed grids (Vazza+17) and SPH (Valdarnini19): at each position, increase $L(\mathbf{x})$ until $\delta \mathbf{v}(\mathbf{x})$ converges.

Extended for AMR, and included in vortex.



Turbulence in galaxy clusters (II: some results)



(Vallés-Pérez, Planelles & Quilis, 2021, MNRAS 504, 510) Global statistic of turbulence: Velocity structure functions:

$$\mathcal{S}_{p}(L) = \langle | \mathbf{v}(\mathbf{x} + L \mathbf{\hat{n}}) - \mathbf{v}(\mathbf{x}) |^{p}
angle_{\mathbf{\hat{n}},\mathbf{x}}$$

We study $[S_p(L)](z)$ at some fixed L.

- Accretion rates tightly correlated (although delayed) with velocity fluctuations
- Solenoidal motions are dominant at all scales

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Turbulence in galaxy clusters (II: some results)

• Enstrophy:

$$\epsilon \equiv \frac{1}{2} \boldsymbol{\omega}^2, \qquad \boldsymbol{\omega} \equiv \nabla \times \mathbf{v}$$

Easy evolution equation

$$\frac{\partial \epsilon}{\partial t} + \frac{1}{a} \nabla \cdot (\epsilon \mathbf{v}) = -2H\epsilon$$
$$-\frac{1}{a} \epsilon (\nabla \cdot \mathbf{v}) + \frac{1}{a} \omega (\omega \cdot \nabla) \mathbf{v}$$
$$+ \frac{\omega}{\rho^2 a} (\nabla \rho \times \nabla P)$$

(Vallés-Pérez, Planelles & Quilis, 2021, MNRAS 504, 510) Compression













Hydrodynamical processes in galaxy clusters

(dɛ/dt)_{source} (Gyr⁻³) 19th October 2021 16 / 20

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What does the future hold?

- Systematic observation of clusters outskirts
 - In radio (e.g., SKA)
 - * Accretion shocks of clusters
 - * Particle acceleration
 - ★ Magnetic fields
 - In X-ray (e.g., Athena)
 - * Explore the thermodynamic conditions in the outskirts
 - * How does gas fall into the cluster?
- Large SZ surveys \implies Discovery of high-redshift systems
 - kSZ ⇒ Directly measure gas proper motions!
- Still a lot of work to do to understand the physical processes acting on galaxy and GC formation!

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The complexity of cosmic flows around cosmic voids

- Huge ($R_{\rm eq} \gtrsim 10 \, {\rm Mpc}$), underdense ($\langle \rho \rangle / \rho_B \sim 0.1$) regions filling up most of the Universe
- Notable applications to:
 - Cosmology
 - Galaxy evolution
- Dominated by outflows



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(Ricciardelli+13)

Void Replenishment: How Voids *Accrete* Matter Over Cosmic History



(Vallés-Pérez, Quilis & Planelles, 2021, ApJ Letters 920, L2)

- We defined a sample of ~ 200 ellipsoidal voids, from a $100 \ h^{-1} \ {\rm Mpc}$ simulation
- Using a pseudo-Lagrangian algorithm, we compute mass fluxes through their boundaries
 - 10% of the mass in voids accreted from overdense regions
 - Reaching 35% for a significant fraction
 - Part of the gas in voids could have been partially pre-processed outside

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Thanks! Time for questions & comments?

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Hydrodynamical processes in galaxy clusters

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